

Ion Acceleration: TNSA and beyond

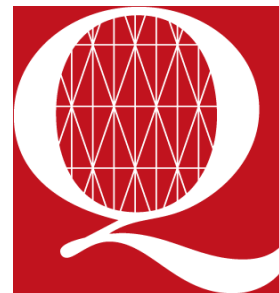
Lecture 1

M. Borghesi

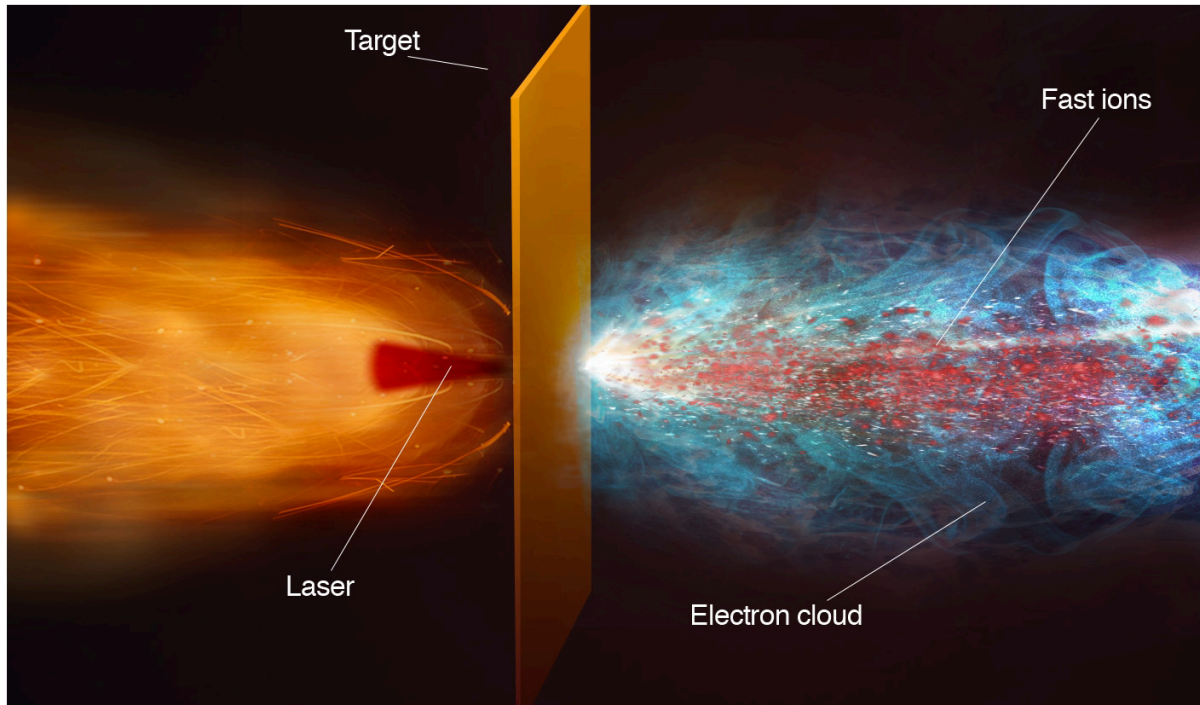
Centre for Plasma Physics,
School of Mathematics and Physics

The Queen's University of Belfast

Advanced Summer School on
laser-driven sources of high
energy particles and radiation
Anacapri, 10-16 July 2017



Ion acceleration : some general points



A.Macchi, M.Borghesi and M. Passoni, *Ion acceleration by superintense laser-plasma interaction*, Rev. Mod. Phys., **85**, 751 (2013)

- Several, fundamentally different mechanisms
- Large accelerating fields sustained by electron-ion separation in a plasma
- Very large fields (up to 10^{13} V/m) applied over very short distances ($\sim\mu\text{m}$)
- Mostly solids (high density targets)

Two classes of lasers are mainly used for ion acceleration

High energy CPA systems

- Nd: Glass technology
- 100s J energy, up to PW power
- Low repetition rate
- 100s fs duration

$$I_{\max} \sim 10^{21} \text{ Wcm}^2$$

VULCAN, RAL (UK)
Phelix, GSI (De)
Trident, LANL (US)
Texas PW, Austin (US)
.....



Ultrashort CPA systems

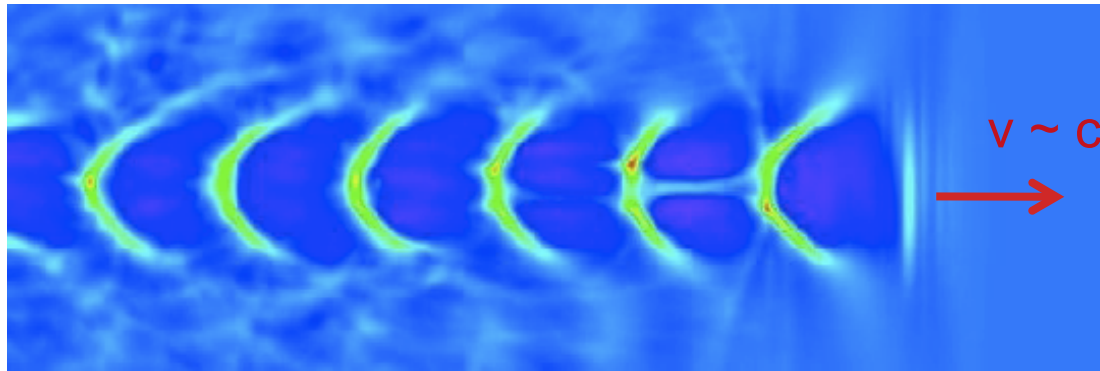
- Ti:Sa technology
- 10s J energy, up to PW power
- 1-10 Hz repetition
- 10s fs duration

$$I_{\max} \sim 10^{21} \text{ Wcm}^2$$

GEMINI, RAL (UK)
Draco, HZDR (De)
Pulser I, APRI (Kr)
J-Karen, JAEA (J)
.....

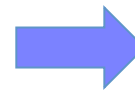


Laser-ion acceleration: why not use wakefield?



Standard laser wakefield methods cannot be used for accelerating ions

- Issues:**
- 1) $v_{\text{ion}} \ll v_{\text{wake}}$
 - 2) v_{ion} varies as ion is accelerated



Need of “slow wave”

Need of accelerating structure with variable phase

Some ideas:

S. Masuda, T. Katsouleas and A. Ogata,
Nucl. Instr. and Meth. A **455** (2000) 172 - 175.

F. Peano, et al, New J. Phys., **10**, 033028 (2008)

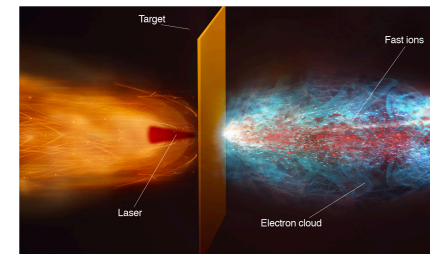
Slow waves from Raman backscattering

Beat wave structure with variable phase velocity (frequency chirped pulses)

Breakdown of lectures

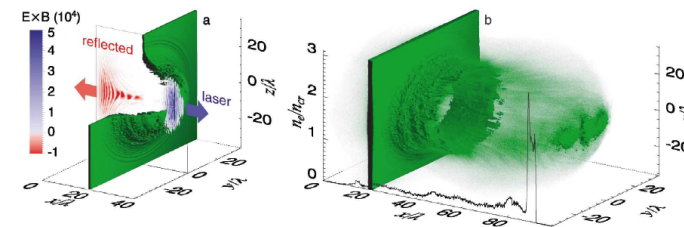
- Lecture 1 : **Sheath acceleration processes**

(Tue, 9 am , 40 min)



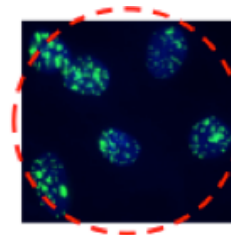
- Lecture 2: **Other mechanisms - new developments**

(Thu, 3pm, 50 minutes)



- Lecture 3: **Applications**

(Friday, 6.30 pm, 30 minutes)



Outline of Lecture 1

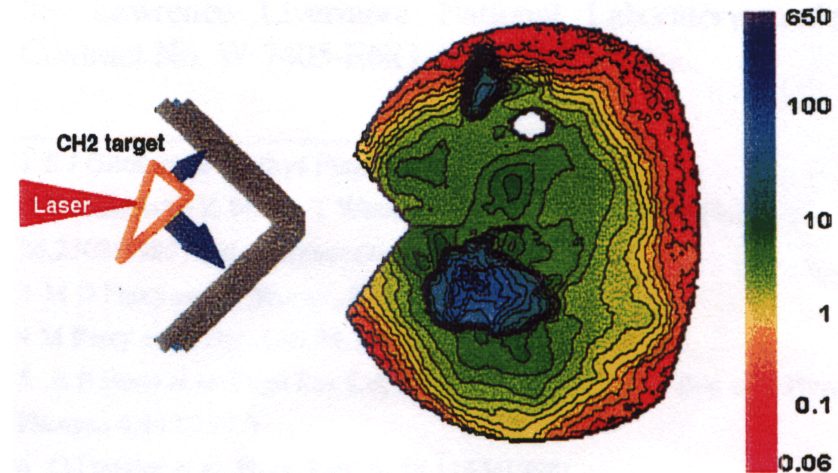
- **Historical introduction**

First observations

- **Target Normal Sheath Acceleration**

The basic process

State of the art and beam properties

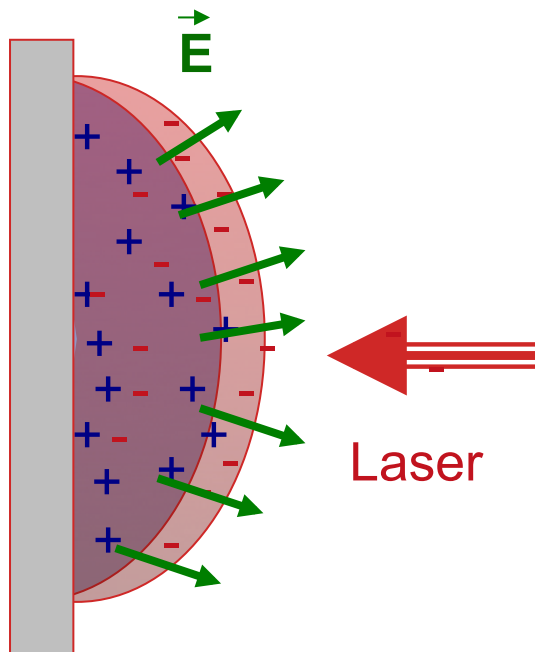


Laser acceleration of ions from laser irradiated targets

was studied from 1960s throughout the 90s

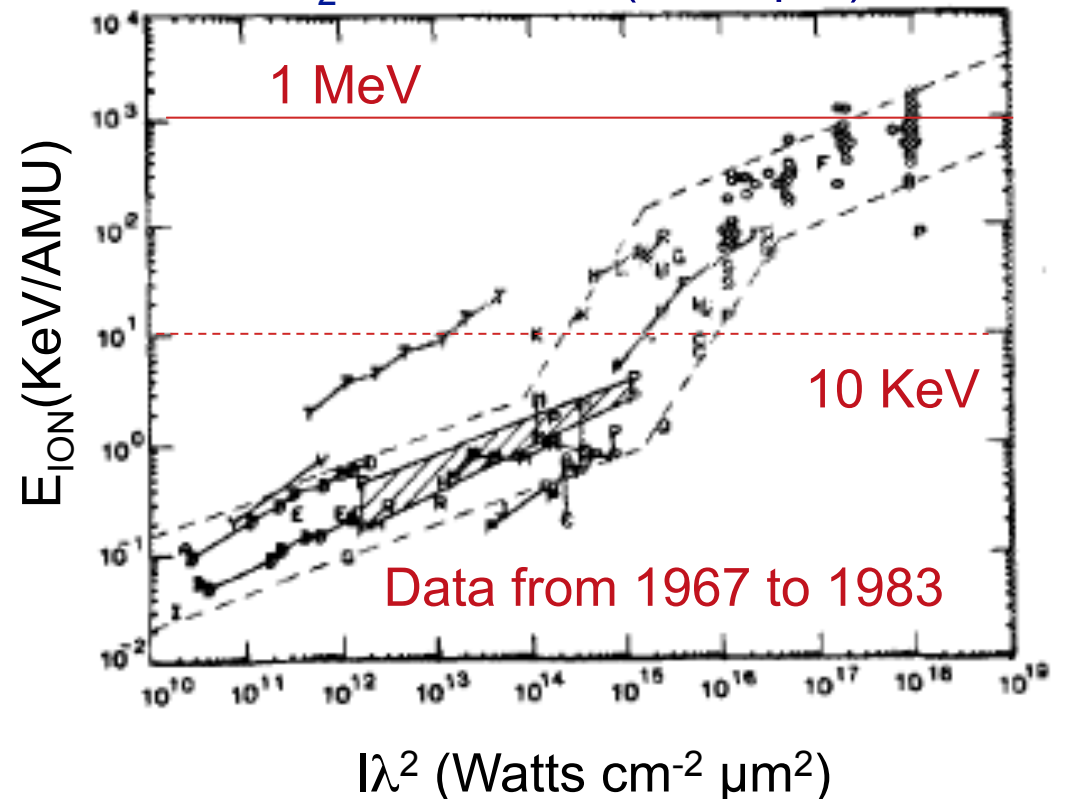
Laser couples energy into electrons

Faster electrons drag ions



A.V. Gurevich et al,
Sov. Phys. JETP, **22**, 449 (1966)

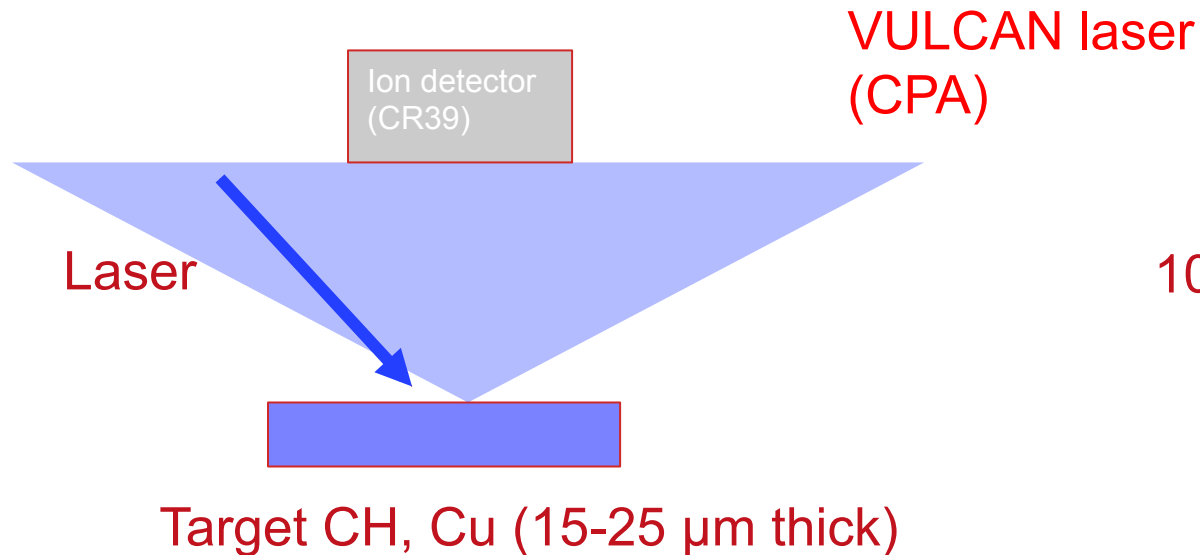
CO₂ lasers , ns, ($\lambda \sim 10 \mu\text{m}$)



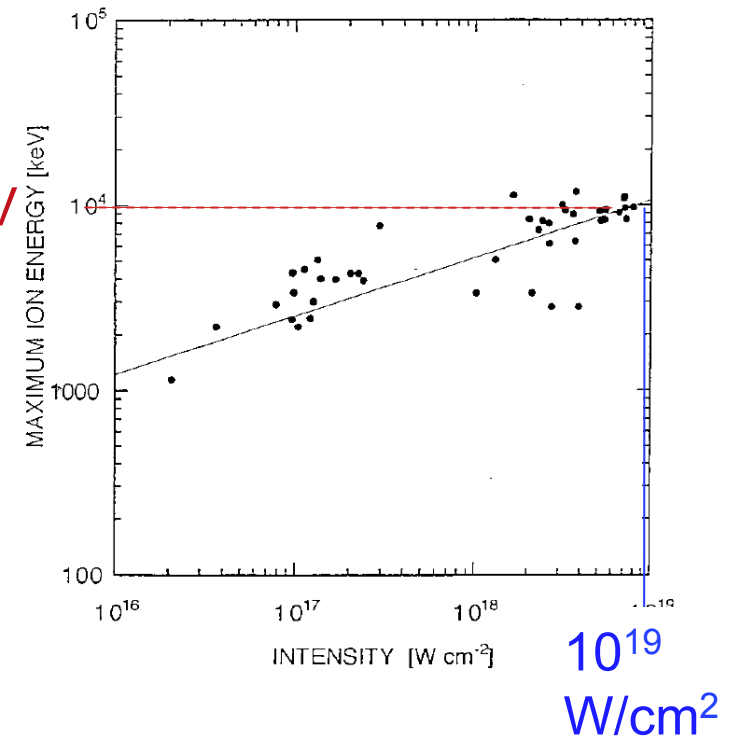
S.J. Gitomer et al,
Phys. Plasmas, **29**, 2679 (1986)

Ion acceleration was studied throughout the 90s

F. Beg et al, Phys. Plasmas, 4, 447 (1997)



10 MeV

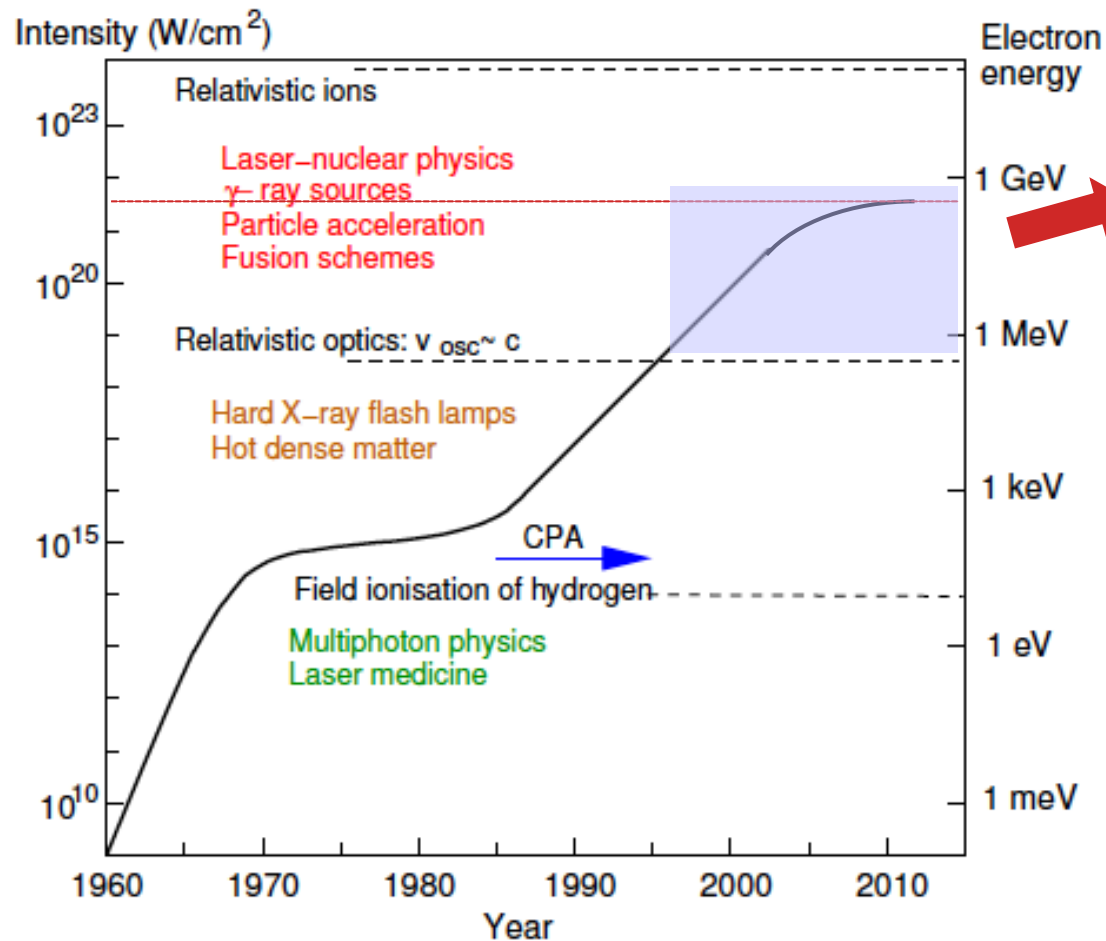


- Protons originate from contaminants
- Exponential spectrum with sharp cut-off
- Ions accelerated during plasma expansion (isothermal rarefaction model)
- Maximum energy of ions related to T_{hot}

$$E_{\text{max}} \sim I^{0.3}$$

(Beg's law)

Laser intensity is key to efficient particle acceleration



Chirped Pulse Amplification (CPA) Techniques

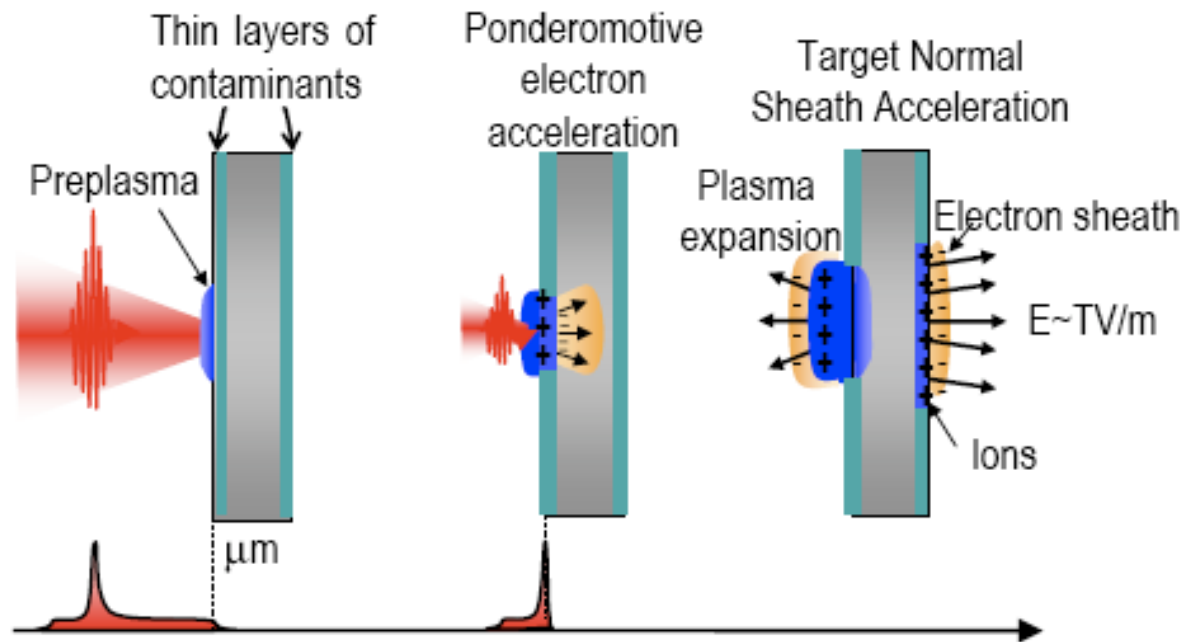
Short, high energy laser pulses



Relativistic interaction regimes
($v_{osc} \sim c$)

Efficient and directional coupling of laser energy into energetic particles

Target Normal Sheath Acceleration (TNSA) from the rear of thin foils was studied from ~ 2000



Clark et al, PRL,
84, 670 (2000)

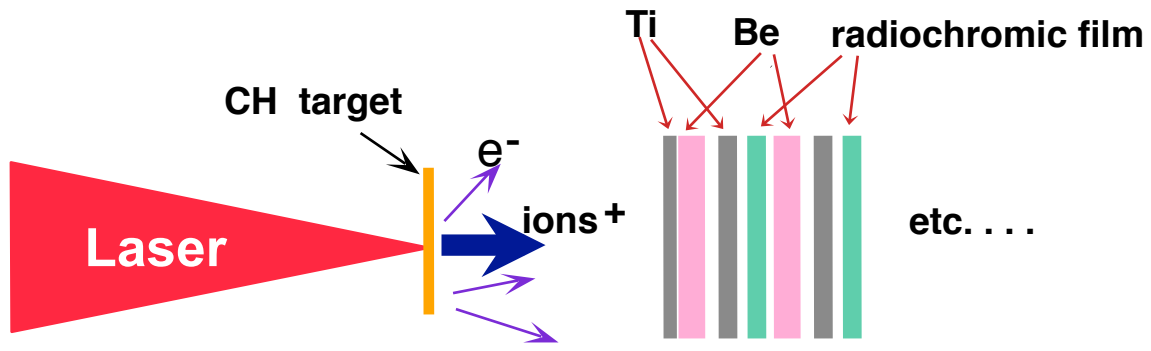
Maksimchuk et al, PRL,
84, 4108 (2000)

Snively et al, PRL,
85, 2945 (2000)

Intensities rising above 10^{19} W/cm^2 – electron acceleration to MeV energies
Thin foils allow electrons to reach the rear of the target and establish a field there
Protons (from contaminants) have *beam* features contrary to lower energy, isotropic emission previously observed from the front.

LLNL data were obtained with Radiochromic film techniques

R. Snavely et al., PRL, 67, 84 (2000)



Counting activations

$Ti^{48} (p,n) V^{48}$ $t_{1/2}=16d, E_{thres}=5 MeV$

gives absolute #'s of protons in several spectral bands.

Spatial distribution of activation gives distribution of protons.

Radiochromic film layers:

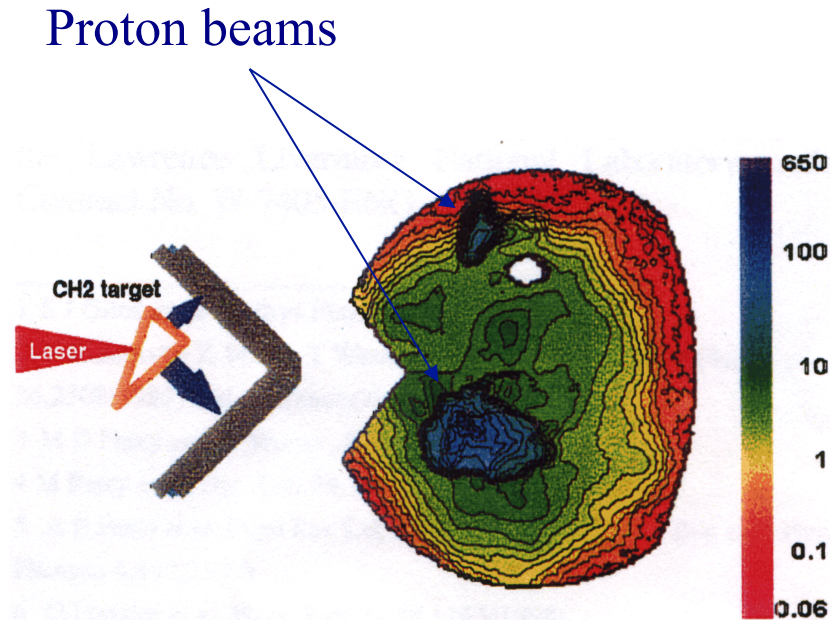
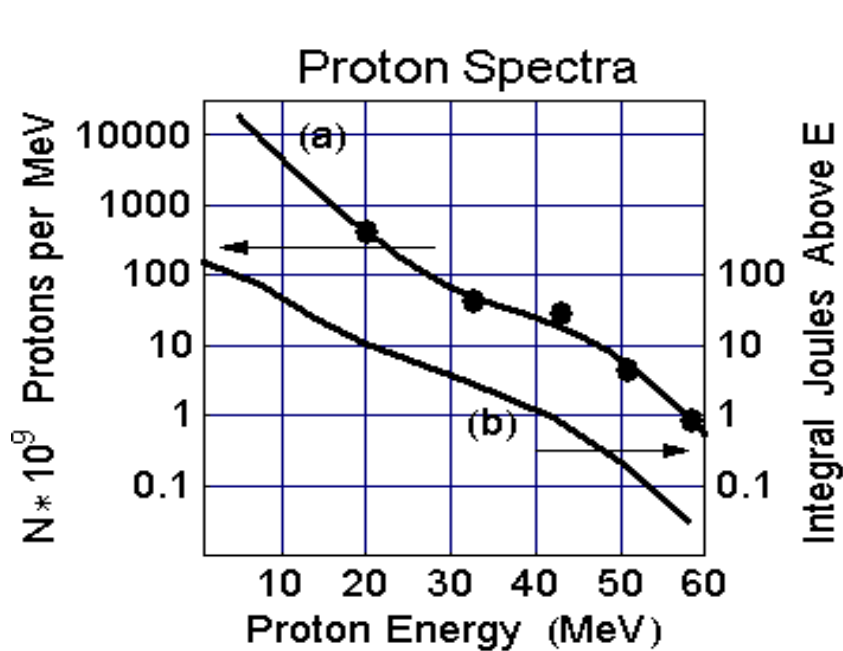


Autoradiographs of activated d Ti discs:



Lawrence Livermore National laboratory experiments

R.D.Snavely *et al* , Phys. Rev. Lett., **85**, 2945 (2000)



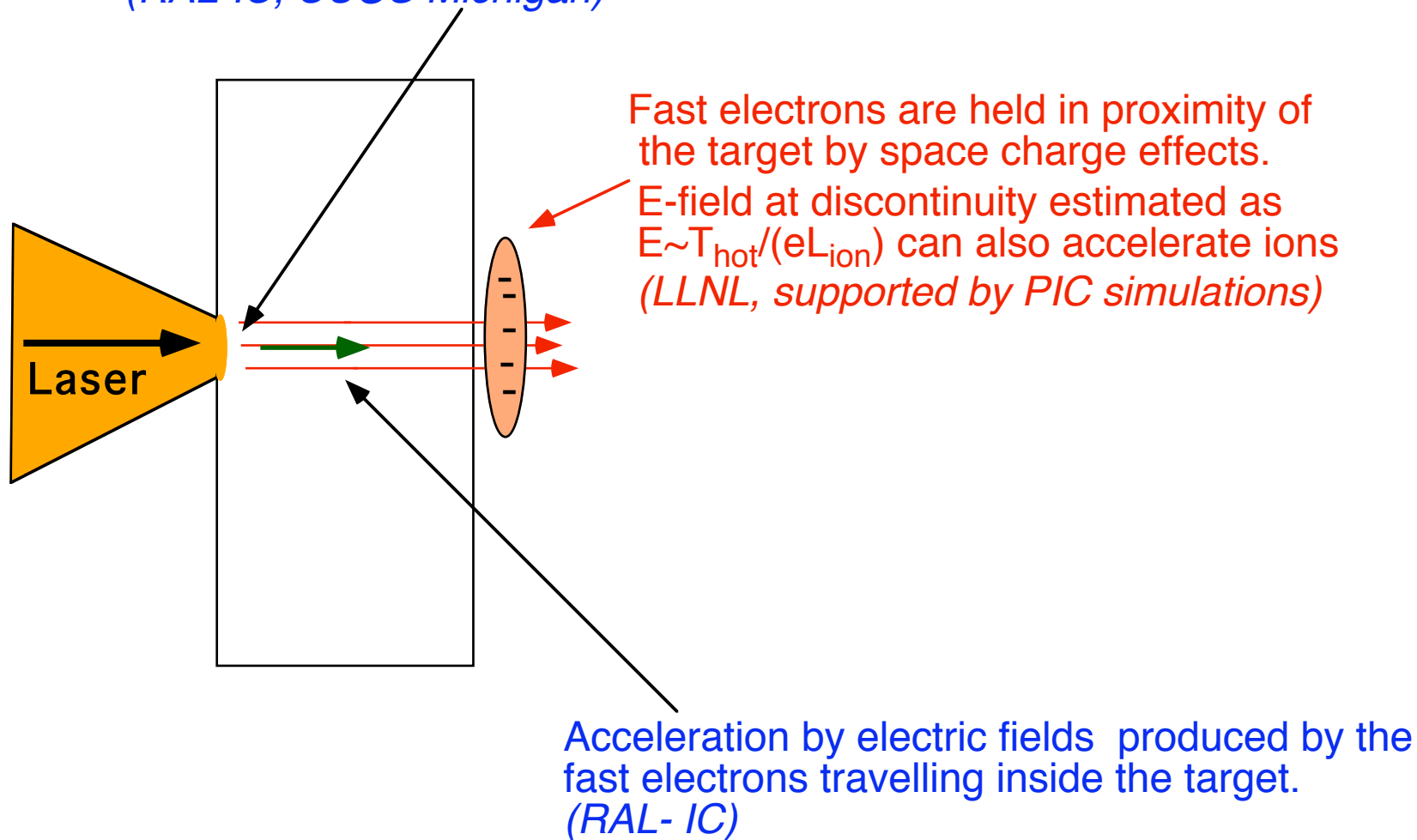
Radiochromic Film (Rads)

- Sharp edge proton beam with cut off energy of 50MeV and slope of 6MeV
- Higher energy protons were more collimated than lower energy
- The proton beam was always normal to the back surface of the target
- No protons > 8 MeV were observed from the front of the target

PW beam
 500J, 0.5ps
 $I \sim 4 \times 10^{20} \text{Wcm}^{-2}$

In 2000 there was controversy regarding the explanation of results

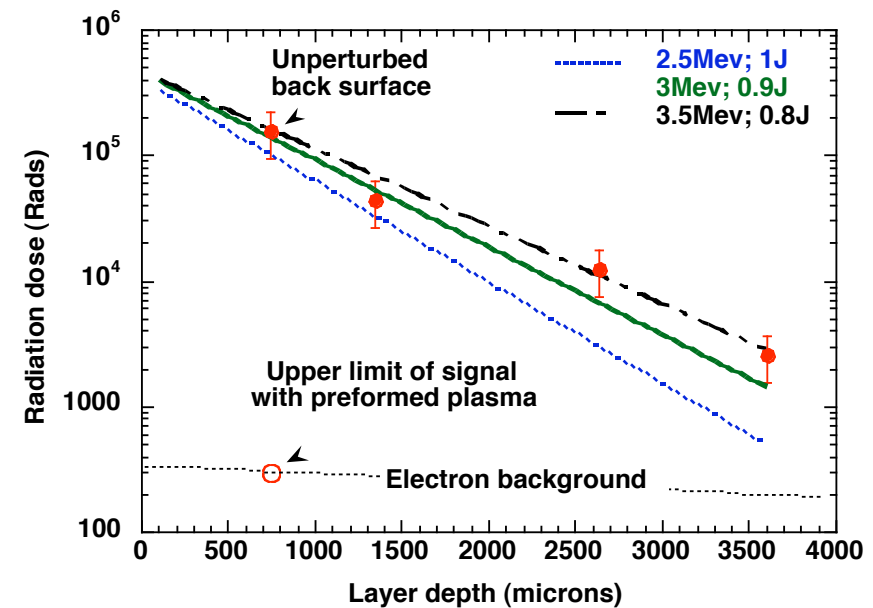
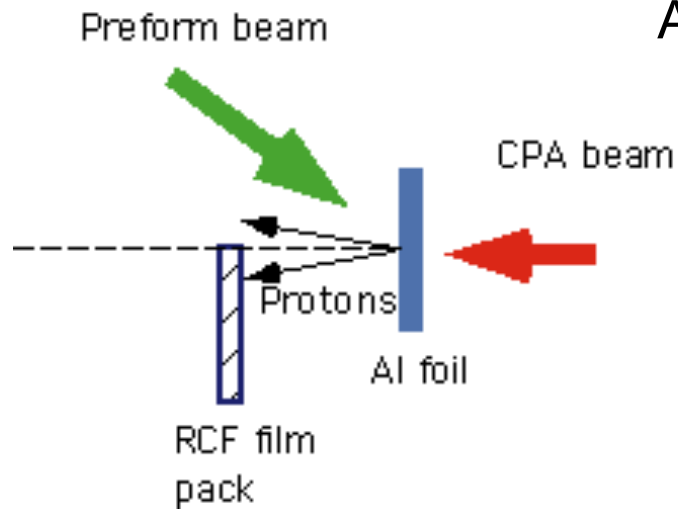
Ponderomotive acceleration (up to a few MeV) followed by electrostatic acceleration due to electron escape
(RAL-IC, CUOS Michigan)



Several experiments have since shown that the dominant acceleration process for “thick” targets takes place at the target rear

Creating a gradient artificially on the rear surface changes very significantly the proton beam energy

A.J.Mackinnon *et al*, Phys. Rev. Lett, 86, 1769 (2001)



Also:

J.Fuchs et al, Phys. Rev. Lett., 99, 015002 (2007)

J.Fuchs et al, Phys. Rev. Lett, 94, 045004 (2005).

Lecture 1 - TNSA

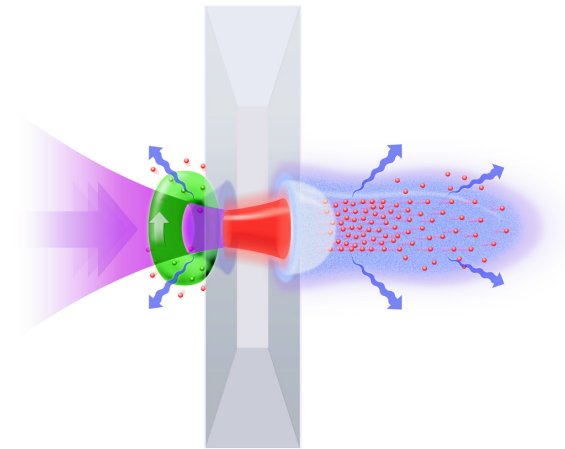
- **Historical introduction**

First observations

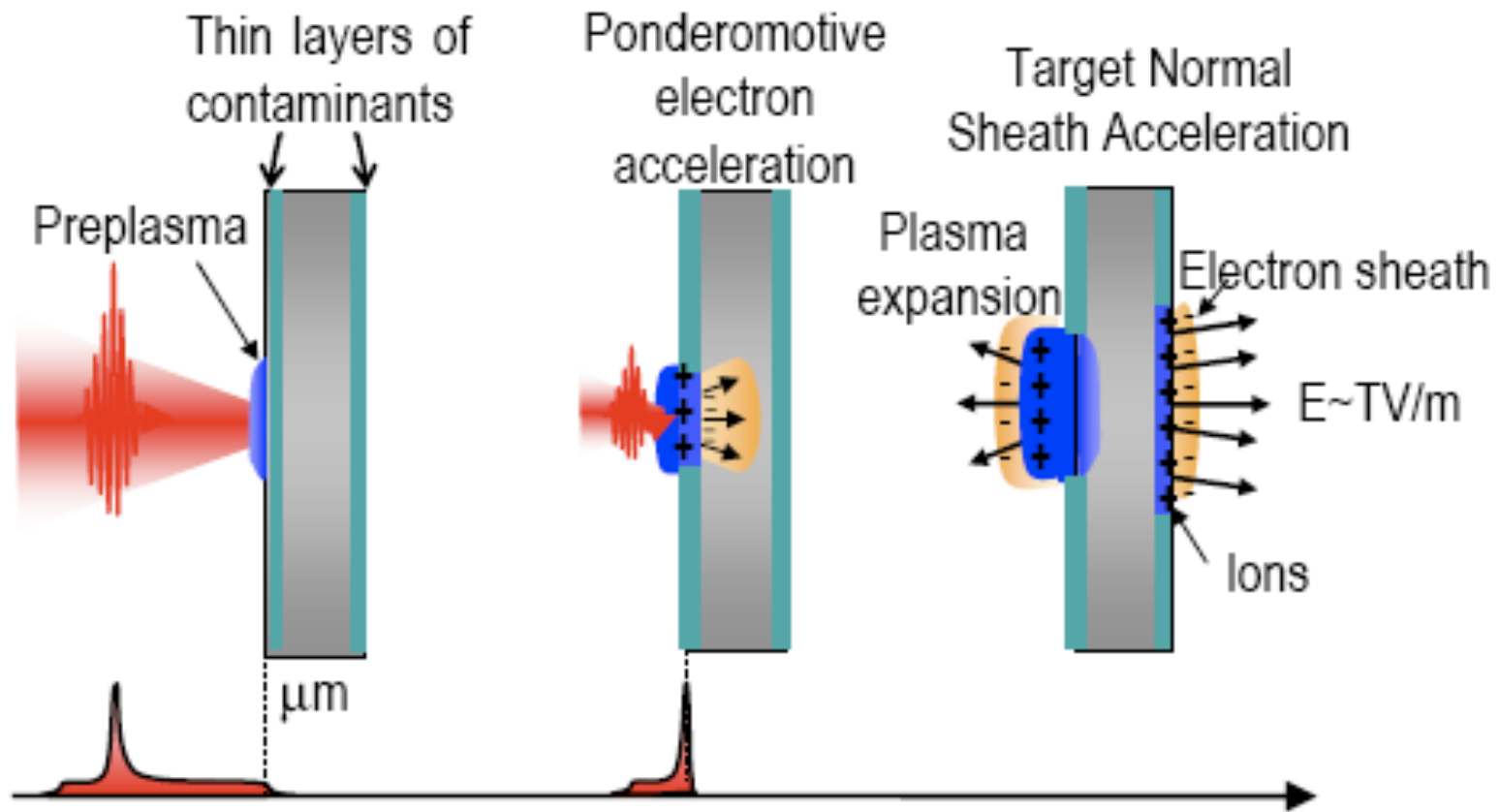
- **Target Normal Sheath Acceleration**

The basic process

State of the art and beam properties

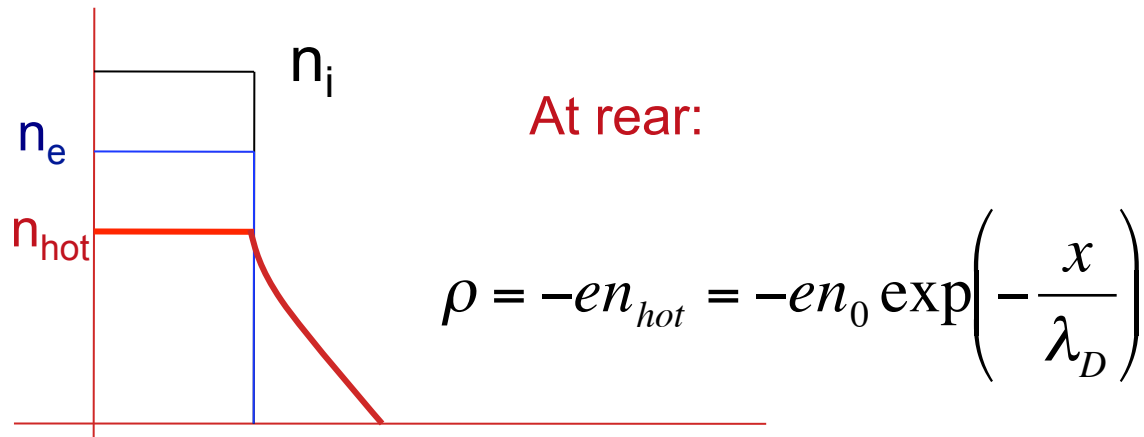


The established mechanism: Sheath Acceleration (TNSA)



High density, high energy electrons lead to ultralarge field

S. Wilks *et al*, PoP, 8, 542 (2001)



$$\nabla \cdot E = \frac{dE}{dx} = \frac{\rho}{\epsilon_0}$$

$$E(0) = \int_{\infty}^0 \frac{\rho}{\epsilon_0} dx$$

$$E(0) = \frac{KT_h}{e\lambda_D} = \sqrt{\frac{n_h KT_h}{\epsilon_0}}$$

Typical values:

$$\lambda_D \sim 1 \mu\text{m}$$

$$T_h \sim \text{MeV}$$



$$E(0) = \frac{10^6 \text{ V}}{10^{-6} \text{ m}} \sim \text{TV} / \text{m}$$

Conventional particle accelerators use much smaller fields



Acceleration by **much smaller** Electric fields associated to alternating voltages (at RF or microwave frequencies)

$$E_{\max} \sim 50 \text{ MV/m}$$

(more than 10,000 smaller than with lasers)



TNSA ion beam properties

- **Low emittance/ high laminarity**
- **Ultrashort duration** (~ ps at the source)
- **High brightness:** 10^{11} – 10^{13} protons/ions in a single shot (> 3 MeV)
- **High current** (if stripped of electrons): kA range
- Divergent (~ 10s degrees)
- Broad spectrum

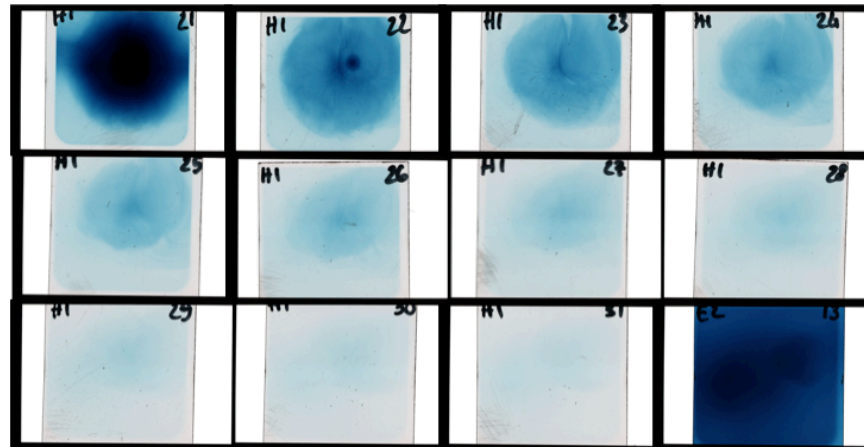
Very compact: $E \sim 1-10$ TV/m
Acceleration lengths: $\sim \mu\text{m}$

Ion beam from TARANIS facility, QUB

$E \sim 10$ J on target in $10 \mu\text{m}$ spot

Intensity: $\sim 10^{19}$ W/cm², duration : 500 fs

Target: Al foil 10um thickness



Laser driven beams have excellent emission quality

Highly laminar source (virtual point source of $\sim\mu\text{m}$ size \ll real source)

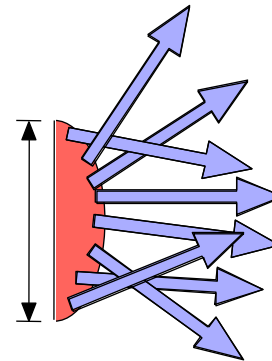
Ultralow emittance/ virtual source:

$$\varepsilon_N < 0.1 \pi \text{ mm.mrad @ 15 MeV}$$

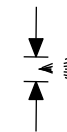
(< 0.004 mm-mrad;

T.Cowan et al, PRL, **92**, 204801, 2004)

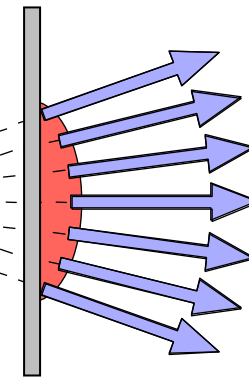
Extended
"thermal"
source



Virtual point
source

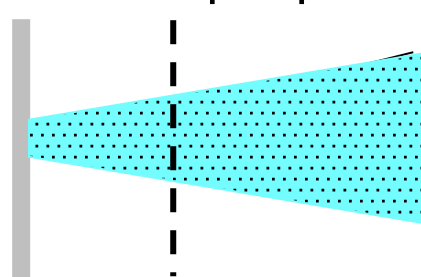


Laminar
source



15 MeV protons

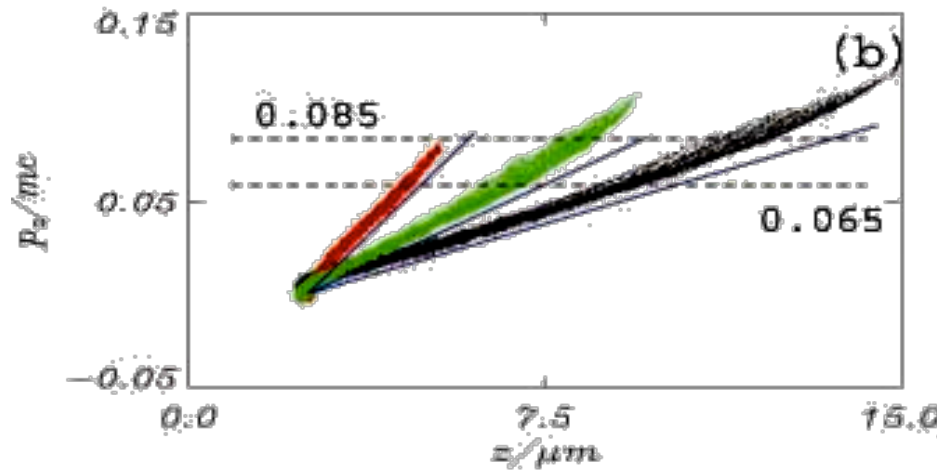
Mesh with $12 \mu\text{m}$ pitch



CERN proton rf-linac:
 $\varepsilon = 1.7 \pi$ mm-mrad

M.Borghesi et al, Phys Rev Lett., 92, 055003 (2004)

PIC simulations predict an excellent longitudinal emittance

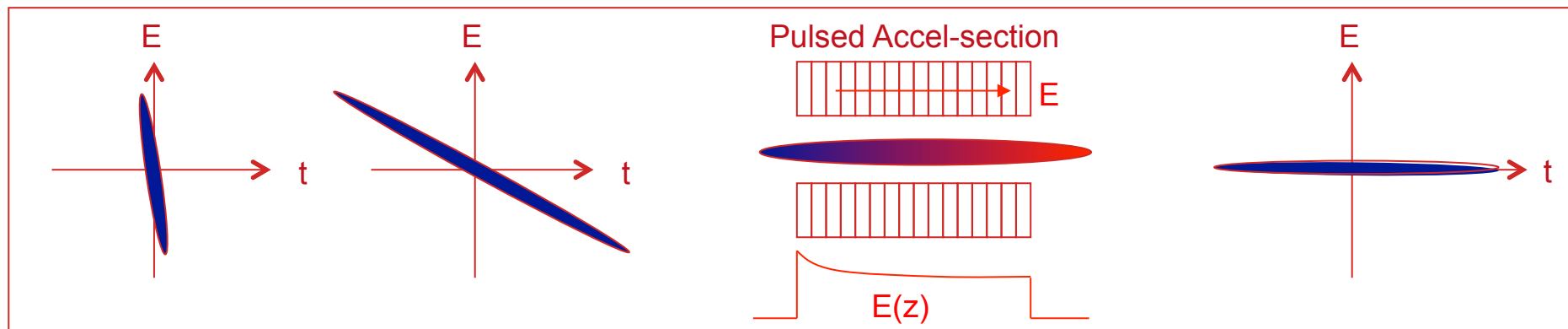


Rapid acceleration produces strong $\Delta E-\Delta t$ correlation

$$\Delta E-\Delta t < 10^{-6} \text{ eV-s}$$



Energy- or time-bunching may be possible with post-acceleration

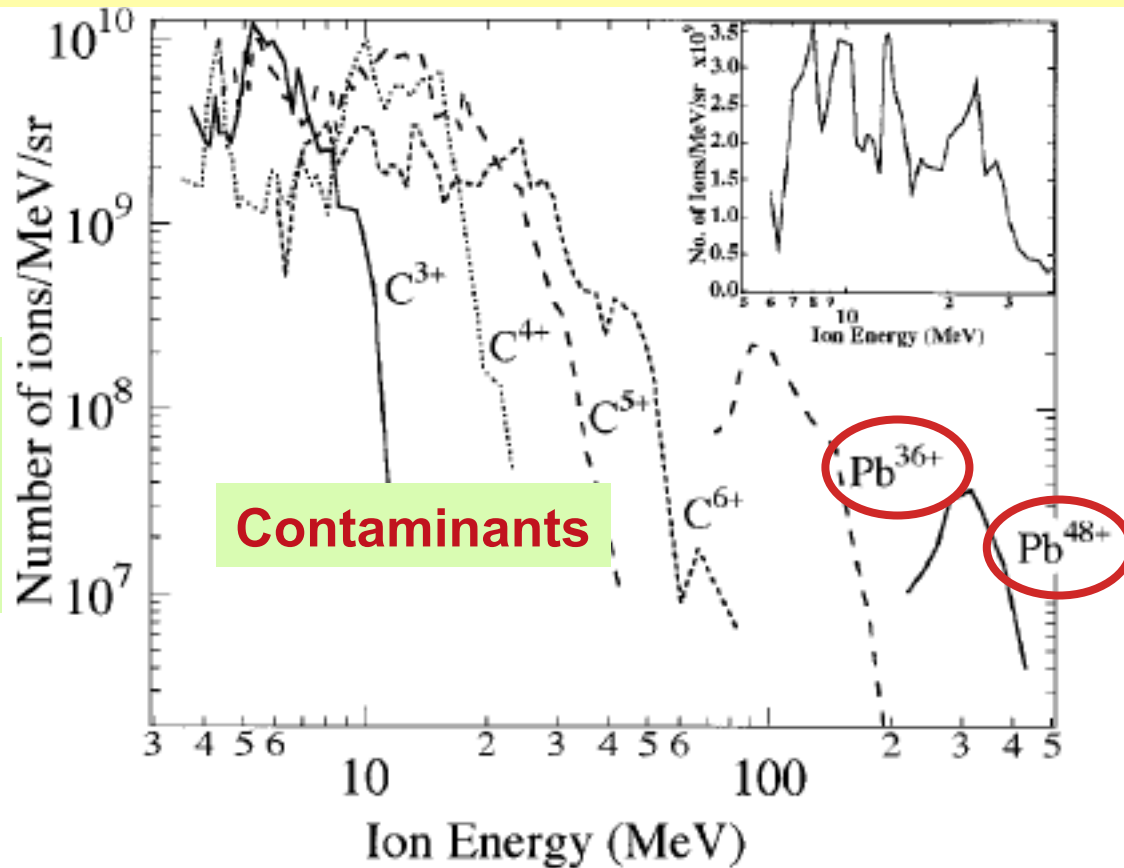


Not only protons but heavier ions also accelerated

Spectra from
VULCAN 100 TW:
Pb target irradiated
@ $5 \cdot 10^{19} \text{ W/cm}^2$

- Most of energy goes to protons
- Energy increases with charge

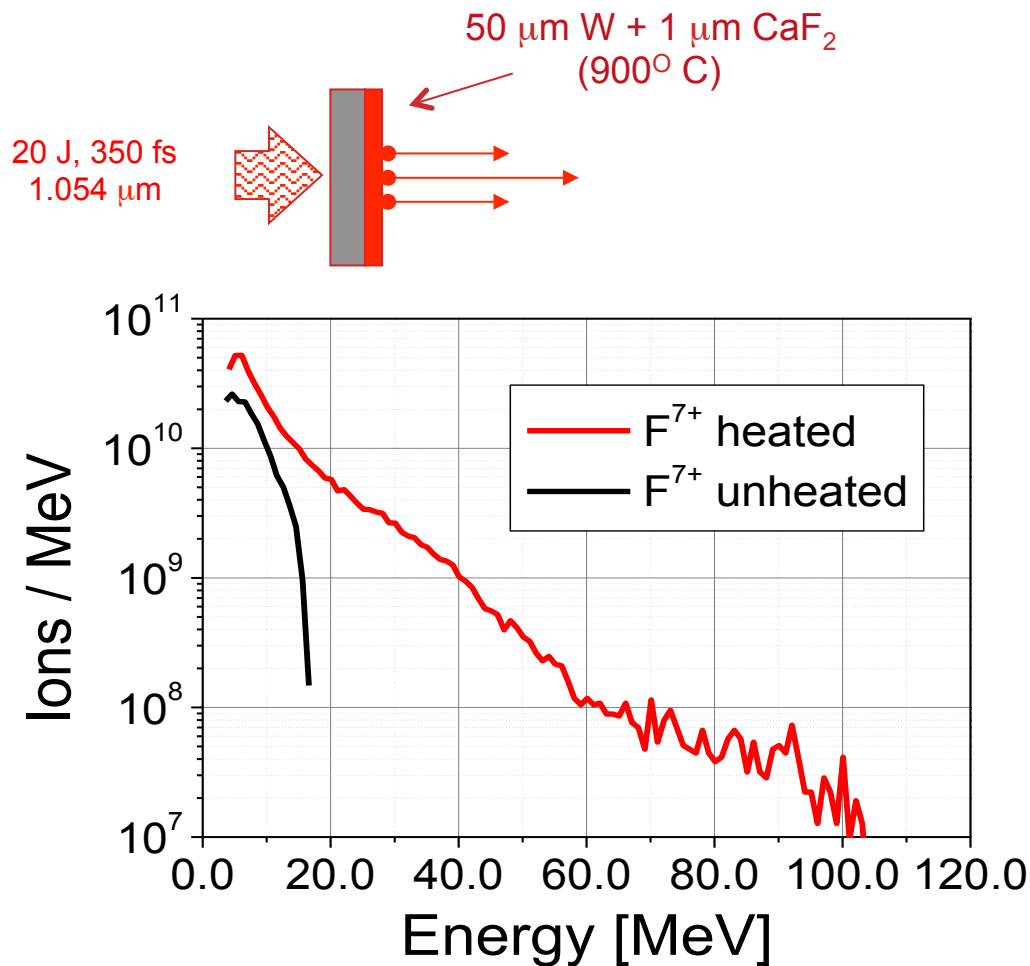
E.L.Clark *et al.*, Phys. Rev. Lett., **85**, 1655 (2000))



Varying the target, any ion can be accelerated (not straightforward with RF accelerators)

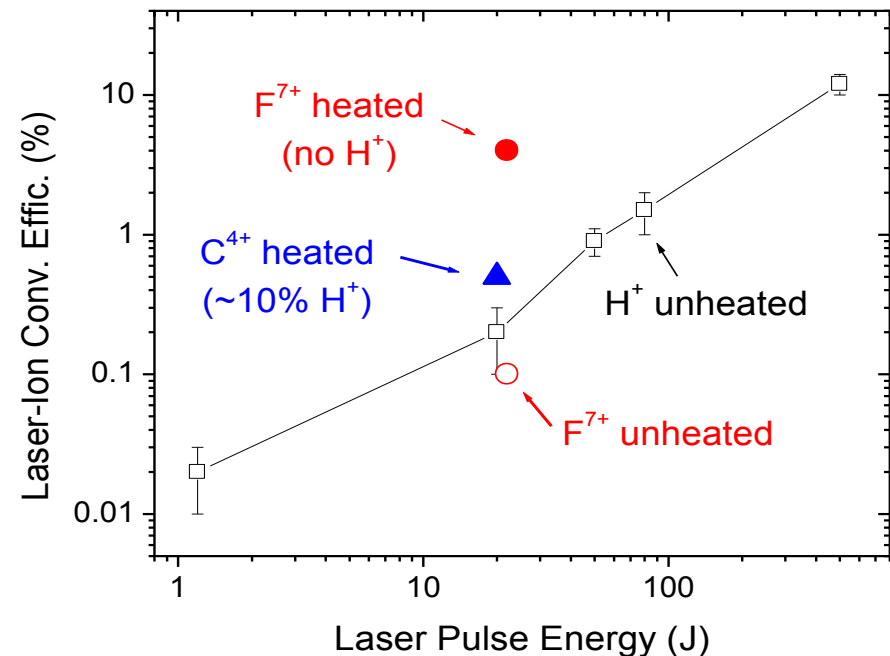
High efficiency conversion of laser energy to heavy ions is achieved by removing hydrogen contaminants from target

M. Hegelich et al., Phys. Rev. Lett. **89**, 085002 (2002)



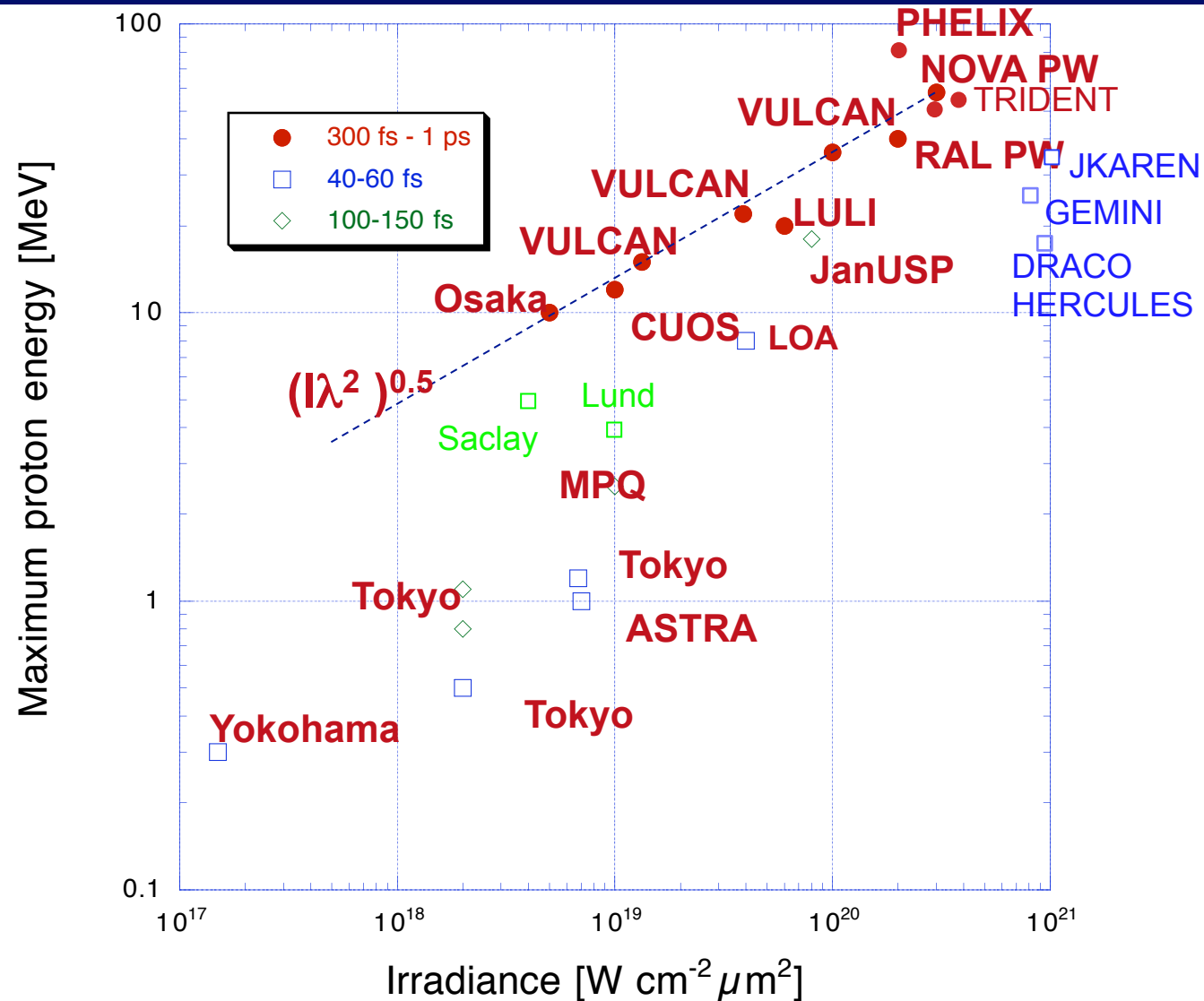
Contaminant removal:

- Ohmic heating
- Laser heating
- Laser ablation



TNSA energies – state of the art

A. Macchi, M. Borghesi, M. Passoni, Rev. Mod. Physics, **85**, 751 (2013)



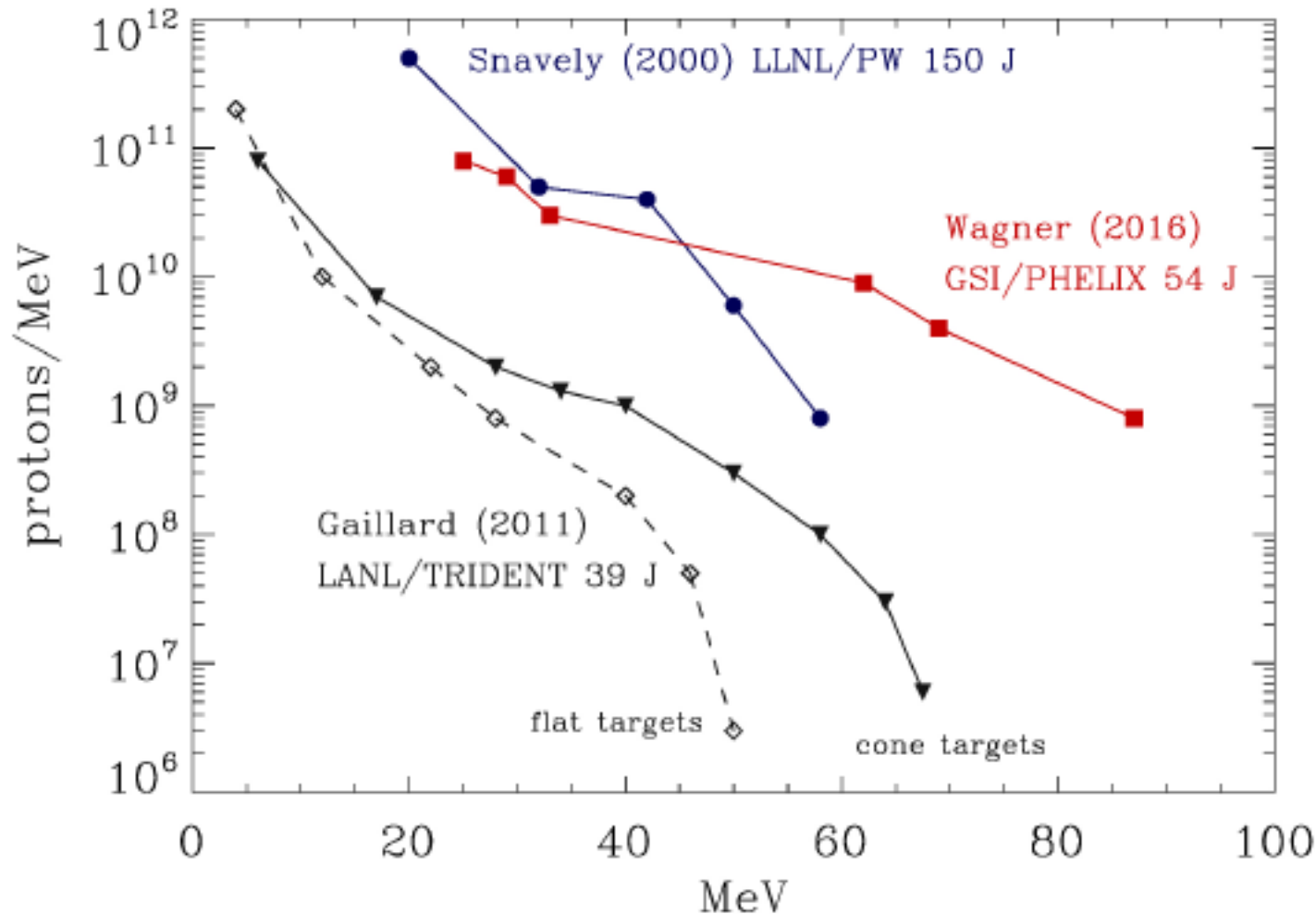
Maximum published energies: ~85 MeV

Conversion efficiency: ~ few %

Acceleration more effective with higher energy, longer pulses, at equal intensities

Effective on protons, less so on higher-Z species

“Record” spectra - long pulses (0.5-1 ps)



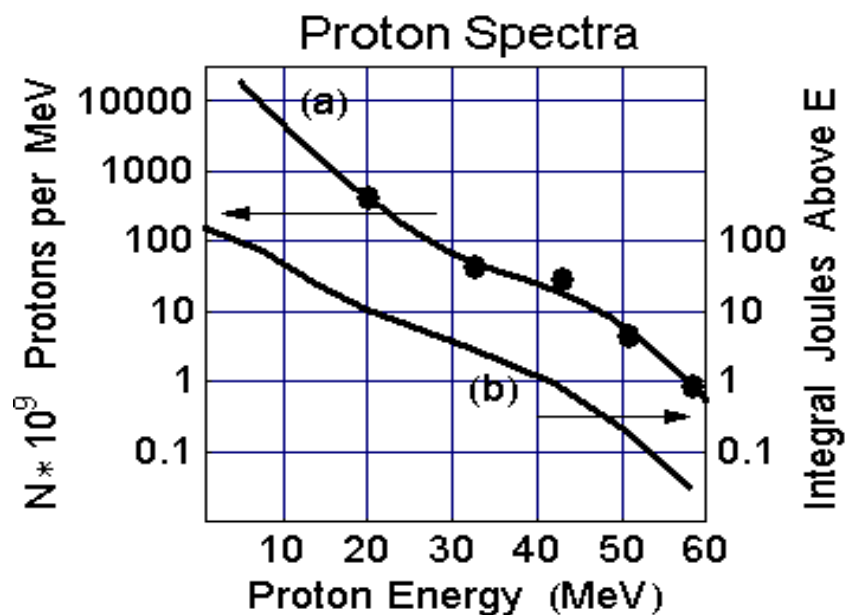
Large, high energy
systems

RCF (and activation)
measurements

$I \sim \text{mid } 10^{20} \text{ W/cm}^2$

Improved control of the laser parameters can lead to significant improvement

R.D.Snavely *et al* , PRL, 85, 2945 (2000)

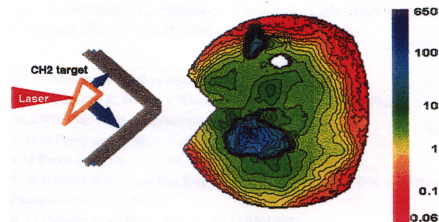


NOVA PW, LLNL

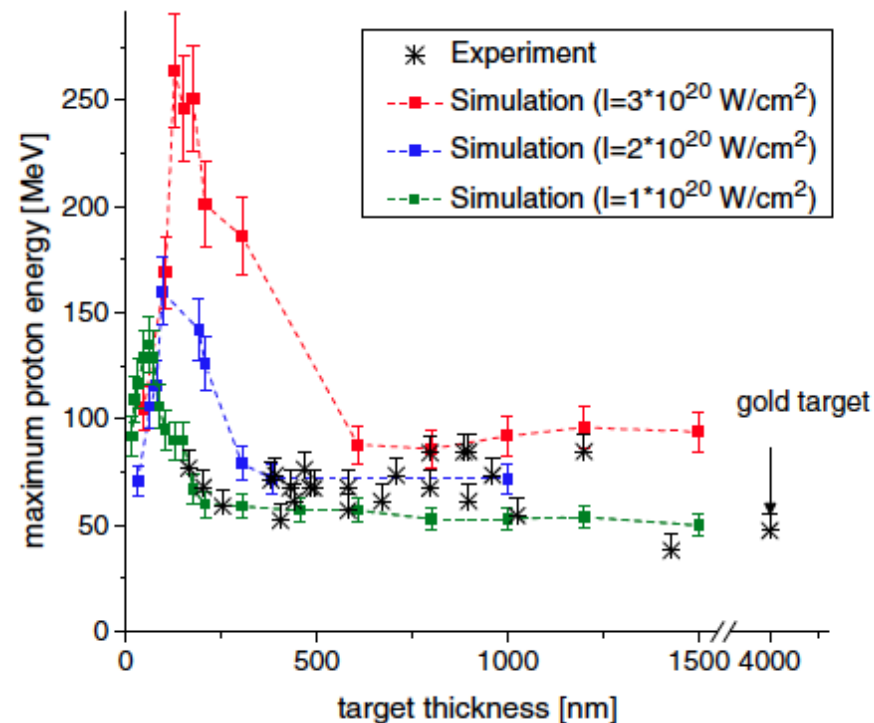
500J, 0.5ps

$I \sim 4 \times 10^{20} \text{ W cm}^{-2}$

CH target, 100 μm thick



F.Wagner *et al* , PRL, 116, 205002 (2016)



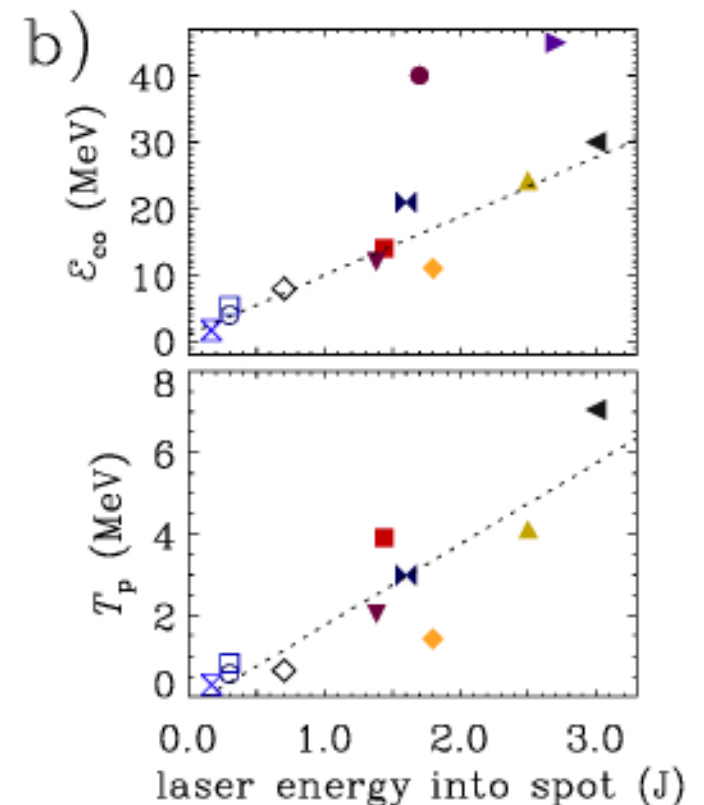
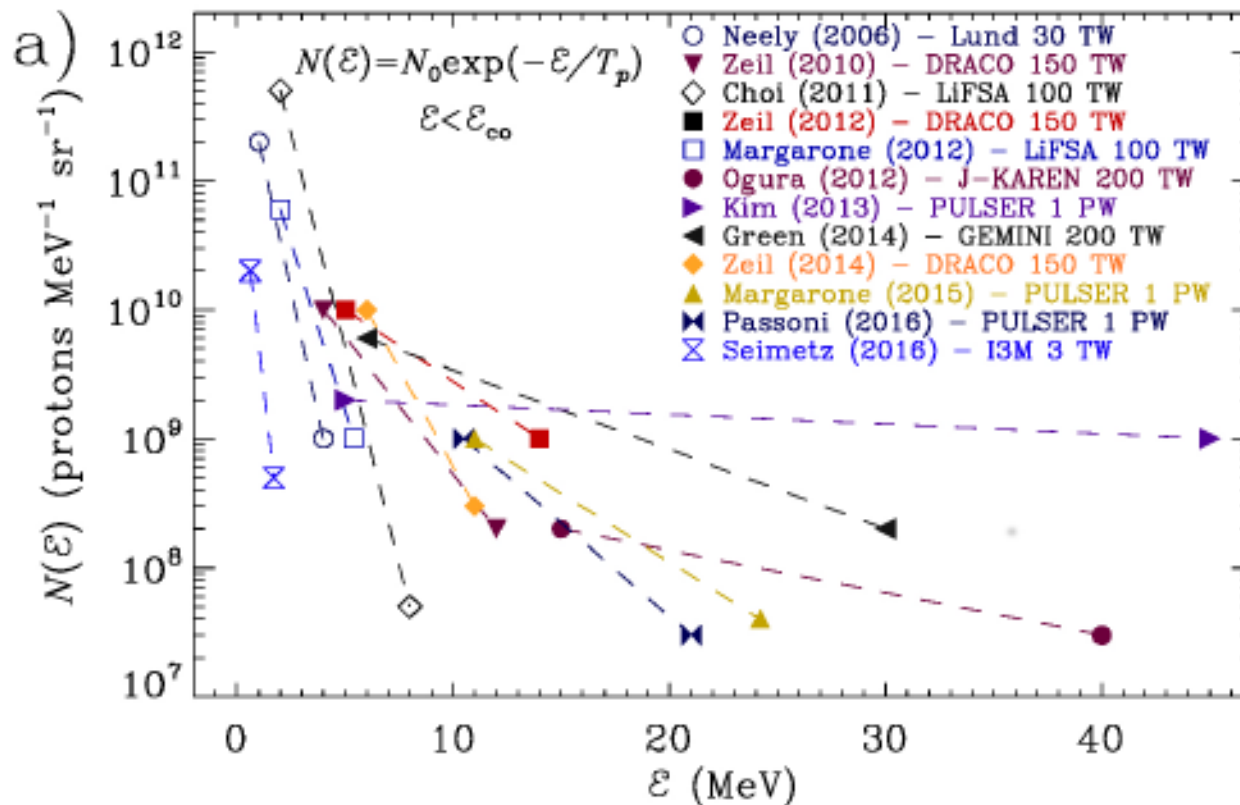
PHELIX, GSI

200J, 0.5 ps, $I \sim 2 \cdot 10^{20} \text{ W/cm}^2$

CH₂ target, 0.9 μm

Better focusing, prepulse control....

Proton spectra from *short pulse* laser systems : ~ 20-50 fs



Near-linear scaling of proton energy with Laser energy/intensity (~ 9 MeV/J)

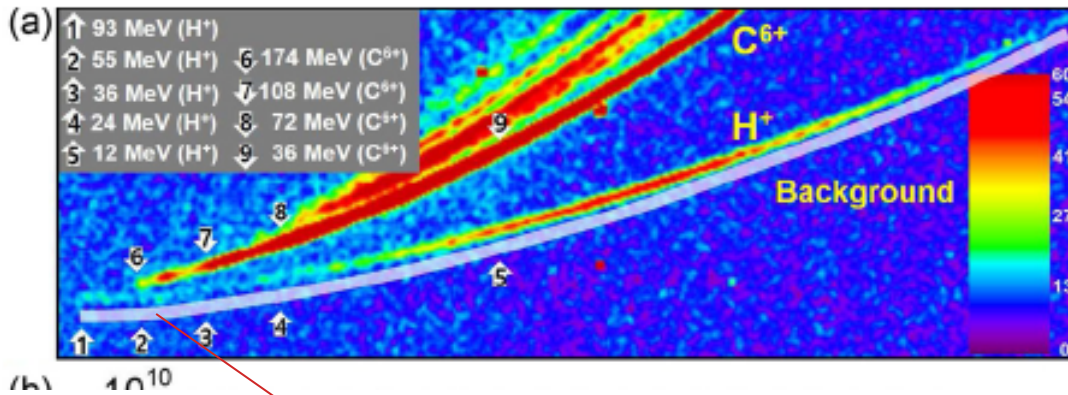
A Macchi, A Sgattoni, S Sinigardi, M Borghesi, and M Passoni, *PPCF*, **55**, 124020 (2013)

Some other high energy claims

PHYSICS OF PLASMAS 23, 070701 (2016)

Radiation pressure acceleration of protons to 93 MeV with circularly polarized petawatt laser pulses

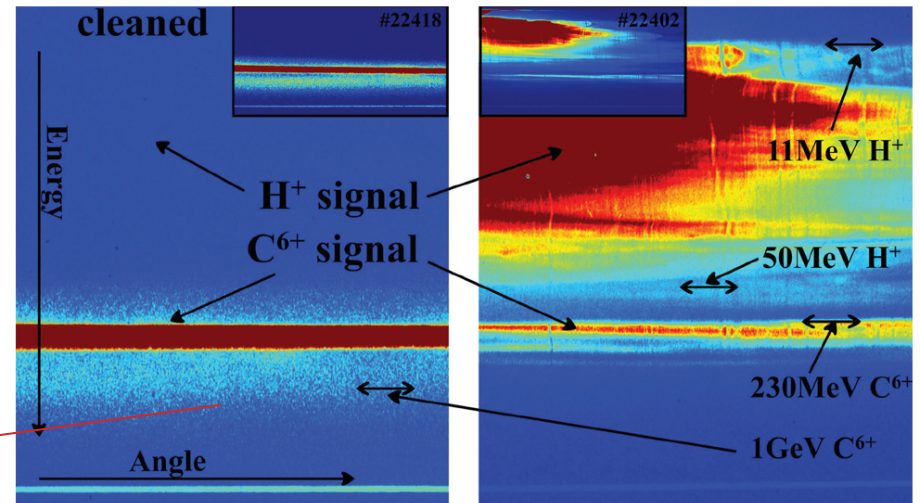
I. Jong Kim,^{1,2,a)} Ki Hong Pae,^{1,2} Il Woo Choi,^{1,2} Chang-Lyoul Lee,² Hyung Taek Kim,^{1,2} Himanshu Singhal,¹ Jae Hee Sung,^{1,2} Seong Ku Lee,^{1,2} Hwang Woon Lee,¹ Peter V. Nickles,³ Tae Moon Jeong,^{1,2} Chul Min Kim,^{1,2,b)} and Chang Hee Nam^{1,4,c)}



PHYSICS OF PLASMAS 20, 083103 (2013)

Laser-driven 1 GeV carbon ions from preheated diamond targets in the break-out afterburner regime

D. Jung,^{1,2,3,a)} L. Yin,¹ D. C. Gautier,¹ H.-C. Wu,¹ S. Letzring,¹ B. Dromey,⁴ R. Shah,¹ S. Palaniyappan,¹ T. Shimada,¹ R. P. Johnson,¹ J. Schreiber,^{2,3} D. Habs,^{2,3} J. C. Fernández,¹ B. M. Hegelich,¹ and B. J. Albright¹



Scattering noise interpreted as signal ?

Need for community established protocols
Rigour in analysis